Resonant x-ray scattering (RXS) has recently become an increasingly important tool for the study of ordering phenomena in correlated electron systems. Yet, the interpretation of RXS experiments remains theoretically challenging because of the complexity of the RXS cross section. Central to this debate is the recent proposal that impurity-induced Friedel oscillations, akin to quasi-particle interference signals observed with a scanning tunneling microscope (STM), can lead to scattering peaks in RXS experiments. The possibility that quasi-particle properties can be probed in RXS measurements opens up a new avenue to study the bulk band structure of materials with the orbital and element selectivity provided by RXS. We test these ideas by combining RXS and STM measurements of the heavy fermion compound CeMn$_2$ ($M = Co, Rh$). Temperature- and doping-dependent RXS measurements at the Ce-$M$ edge show a broad scattering enhancement that correlates with the appearance of heavy $f$-electron bands in these compounds. The scattering enhancement is consistent with the measured quasi-particle interference signal in the STM measurements, indicating that the quasi-particle interference can be probed through the momentum distribution of RXS signals. Overall, our experiments demonstrate new opportunities for studies of correlated electronic systems using the RXS technique.
RESULTS

Our earlier STM studies (32, 33) showed that the cleaved (001)-oriented surfaces of CeCoIn₅ expose three different chemical terminations: surfaces A, B, and C. In these previous works, we demonstrated the surface-dependent sensitivity to the heavy fermion excitations: The light quasi-particles were detected on surface A, whereas surface B predominantly showed the heavy quasi-particles of the hybridized band structure. Here, we carried out spectroscopic measurements in the normal state (T = 10 K) on surface B of CeCo(In₁₋ₓCdₓ)₅ to study the heavy f–quasi-particle interference and compare it with the 4f-sensitive RXS data. Previously (32), the hallmark of heavy-band formation was observed through the temperature-dependent large density of states in the STM spectra (Fig. 1B). Here, we use SI-STM to visualize its energy-momentum structure. Figure 1 (C and D) shows the real-space conductance map (x = 0.15) at specified energies near Eₚ. The Fourier transforms of the conductance maps (Fig. 1, E and F) reveal an enhancement of local density of states modulations along the [H,H] crystallographic direction. The energy-momentum structure of these quasi-particles in the [H,H] direction signals the presence of rapidly dispersive bands as a function of energy (Fig. 1G). Similar results were obtained for the x = 0.0075 Cd-doped sample (Fig. 1H). Therefore, both results reveal the formation of heavy quasi-particle bands near the chemical potential, which are independent of the Cd doping: The quasi-particle interference is unaffected by the underlying ground state at this temperature. Overall, the dispersive nature of the modulations in the STM conductance maps, with the absence of nondispersive features, relates its origin to the quasi-particle interference of heavy f-electrons. The quasi-particle interference that originates from the heavy bands and the hybridization energy scale agree with ARPES measurements of heavy f–quasi-particles close to the Fermi energy (37, 38).
However, from the experimental perspective, because both RXS and STM are momentum-transfer (Q-space) probes, the comparison between them becomes more direct.

We now move to our RXS measurements performed on CeCo(In$_{1-x}$Cd$_x$)$_5$ samples ($x = 0$ and $x = 0.1$) and on CeRhIn$_5$. To enhance the sensitivity of our scattering measurements to the $f$-electron states, we tuned the photon energy near the Ce-$M_{4,5}$ edges, as determined by the x-ray absorption spectrum (XAS). The XAS in Fig. 2A displays peaks due to the photon energy near the Ce-$3d$ edge to be entirely of $f$-electron states, with a characteristic width of 17 eV, which is the spin-orbit splitting of the $J = 3/2$ and $J = 5/2$ states of the $3d$ core hole, and corresponding to the $M_{4,5}$ edges, respectively (39, 40). Recent dynamical mean field theory calculations (41, 42) show the $f$-quasi-particle peak near the Fermi energy (also seen in our STM data in Fig. 1) to be entirely of $J = 5/2$ character, whereas the $J = 7/2 f$-band is located about 280 meV above the chemical potential. These strongly dispersing $f$-bands near the Fermi energy have been experimentally observed by ARPES measurements (37, 38) and STM measurements (33, 34). Dipole selection rules ($\Delta J = 0, \pm 1$) dictate that whereas at the $M_4$ edge both $J = 5/2$ and $J = 7/2$ unoccupied $4f$ states can be reached, only the $J = 5/2$ state can be reached at the $M_4$ edge. Therefore, we conclude that the RXS measurement at the $M_4$ edge is selectively sensitive to the narrow heavy $f$-quasi-particle peak at $E_F$ and thus allows a direct comparison to the SI-STM data of Fig. 1 in an energy window smaller than the XAS broadening. In the following, therefore, we focus on RXS measurements at the Ce-$M_4$ edge.

The RXS measurements were performed using a standard scattering geometry (Fig. 2B) in the $\{HHL\}$ plane with $\sigma$ polarization and using a photon energy of $E_{\text{ph}} = 900.3$ eV in resonance with the Ce-$M_4$ edge, unless otherwise noted. Because of the geometric limitations imposed by this small photon energy, the momentum scans are not restricted to a single value of momentum transfer $L$ along the crystallographic $c$ axis (see the Supplementary Materials), similar to previous works on the cuprates (5–11).

Figure 3A shows momentum scans plotted as a function of in-plane momentum transfer along the two high-symmetry directions $[H,0]$ and $[H,H]$ in CeCoIn$_5$ at the Ce-$M_4$ edge and also at the Co-$L_2$ edge (corresponding to a $2p$-to-$3d$ transition). All scans exhibit a sharp increase (decrease) of intensity for small $H < 0.15$ rlu (large $H > 0.4$ rlu) related to the geometry of the RXS experiment, where we defined rlu as $2\pi/a = 1.36$ Å$^{-1}$, with $a = 4.6$ Å, the tetragonal in-plane lattice constant. The momentum scans reveal a broad scattering enhancement residing on top of a temperature-independent fluorescence background in the $0.2 < H < 0.4$ rlu range, along the $[H,H]$ direction when tuned to the Ce-$M_4$ edge, but there was no enhancement along the $[H,0]$ direction. Similar momentum scans, with the $x$-ray photon energy tuned to the Co-$L_2$ absorption edge, show the absence of the scattering enhancement along $[H,H]$. Furthermore, momentum scans for photon energies finely tuned around the Ce-$M_4$ resonance (Fig. 3B) demonstrate that the scattering enhancement near $H = 0.35$ rlu is resonant with the $M_4$ edge. These results indicate that the scattering enhancement observed at the $M_4$ edge along the $[H,H]$ direction originates from heavy quasi-particles of $f$-character just above $E_F$.

Typically, resonant enhancement in scattering experiments is associated with electronic ordering. In this context, it must be noted that CeCo(In$_{1-x}$Cd$_x$)$_5$ is located close to an antiferromagnetic (AFM) quantum critical point (35, 36), with Cd doping driving the system toward the AFM ground state ($T_N \sim 3$ K at $x = 0.10$). Hence, although all our measurements were performed in the absence of static order, it would be conceivable that at 10 K our RXS measurements, which are energy-integrated and therefore also sensitive to inelastic processes, could be picking up fluctuations of an ordered state. However, this possibility is repudiated by the insensitivity of the RXS scattering enhancement to Cd doping (Fig. 3C), demanding an alternative explanation.

To further investigate the origin of this resonant scattering enhancement, we next consider its temperature and material dependences. Figure 3D displays momentum scans along the $[H,H]$ direction for several temperatures at the Ce-$M_4$ edge. The data reveal a significant temperature dependence, with a rapid suppression of the scattering enhancement up to 100 K and its saturation above. Identical measurements carried out on the isostructural material CeRhIn$_5$, where heavy $f$-quasi-particles are expected to be absent from the Fermi surface at 20 K (32), show no temperature-dependent scattering features in the same temperature window (Fig. 3E). These temperature- and material-dependent RXS measurements provide a direct connection to the STM

Fig. 2. Resonant transition of the heavy $4f$ state. (A) Linearly polarized XAS spectra [measured through the total electron yield (TEY)] on CeCo(In$_{1-x}$Cd$_x$)$_5$ at $x = 0$ (blue) and $x = 0.1$ (green) doping level at 10 K, which demonstrates the transition between the $3d$ and $4f$ states. The curves are shifted vertically for clarity. (B) RXS scattering geometry, with the scattering plane lying at 45° relative to the $a$ and $b$ axes.
measurements (Fig. 1B) and further provide an important finding. The robust presence of the heavy fermion band in CeCo(In\(_1-x\)Cd\(_x\))\(_5\) for \(x = 0.1\) and 0.15 and its absence in CeRhIn\(_5\) from STM and RXS indicate that the AFM ground state, which forms at lower temperatures in both compounds, has different origins, presumably related to the itinerant or localized character of the Ce\(^{3+}\)'s \(f\)-moments. This is an important piece of information in the context of Kondo destruction and quantum criticality in heavy fermions \((43)\) and deserves further investigation.

Figure 3F displays the amplitude of the scattering peak enhancement in RXS obtained at a given temperature as the difference between the area under the momentum scan at that temperature and at 200 K. Comparison to the temperature dependence of the \(f\)-weight from STM spectra shows a good agreement with the RXS results, indicating not only a strong correspondence between the two techniques but also that RXS can be a momentum-resolved probe of the band structure of materials beyond ordering phenomena.

Exactly how this sensitivity to the band structure occurs remains an open question in the field. To illustrate one possible scenario, in which a strongly dispersing (flat) band above the Fermi energy can give rise to a broad scattering enhancement in the RXS measurement, we follow the procedure proposed by Abbamonte et al. (21) and Dalla Torre et al. (22), using the experimentally obtained quasi-particle interference information rather than simulated data. The derived phenomenological picture relates the Fourier transform of STM conductance maps, \( g(q,\omega) \), to the RXS intensity

\[
I_{RXS}(q, E_{ph}) = \left| \int \frac{g(q, \omega')}{E_{ph} - \omega' - E_h + i\Gamma_h} d\omega' \right|^2, \tag{1}
\]

where the integral runs only over unoccupied states, whereas \( E_h \) and \( \Gamma_h \) refer to the core-hole energy and broadening, respectively. At the resonant condition (that is, \( E_{ph} = E_h \)), this integral effectively relates \( I_{RXS} \) to the sum of \( g(q,\omega) \) inside an envelope of width \( \Gamma_h \), which is typically \( \sim 200 \text{ meV} \).

Before applying Eq. 1 to our experimental data, we discuss the approximations and assumptions that it entails. First, we point out that choosing the proper integration boundaries is critical to the result of Eq. 1. Naturally, the limited energy range (few hundreds of millielectron volts) of an STM measurement requires us to introduce a cutoff energy for the upper bound of the integral. However, as we discussed earlier, because of dipole selection rules, the Ce-M4 edge is expected to represent only the \( j = 5/2 \) heavy-electron states near \( E_F \). Therefore, we restrict the integration only up to 20 meV, where we detect the strongly dispersing \( f \)-band, which is also the energy window where the tunneling spectrum exhibits a temperature dependence (Fig. 1B). Extending the integration to higher energies would suppress the strength of the signal associated with the heavy band relative to the background intensity of \( g(q,\omega) \) and would also introduce spurious contributions from bands that may not play a role in the RXS measurement.

Second, it is important to note that the quasi–three-dimensional nature of the band structure of CeCoIn5 is not accounted for in Eq. 1. Both STM and RXS techniques are sensitive to the \( Q_x \) and \( Q_y \) (or \( H \) and \( K \) in reciprocal lattice units) in-plane components of the \( Q = k_f - k_i \) momentum transfer, as well as the \( Q_z \) (or \( L \)) out-of-plane component. Generally, quasi-particle interference maps measured by STM in a material with three-dimensional band structure can be approximated by the weighted average quasi-particle interference over \( k_z \) slices (44), though the exact nature of this sum is not known. On the other hand, the scattering geometry in RXS measurements precisely determines the value of \( L \) (see the Supplementary Materials), but not the initial and final values of \( k_z \). Therefore, because the sensitivity to \( k_z \) may be different in the two techniques, an exact connection between RXS and STM using Eq. 1 may not hold. The calculations below should be interpreted as a qualitative description.

Figure 4 shows the calculated RXS intensity based on our STM conductance maps acquired on surface B along the \([H,H]\) direction. We observe that the flat band indicated by the dashed area in Fig. 4A may not play a role in the RXS measurement.

**Fig. 4. Quasi-particle interference and RXS.** (A) Energy-momentum structure of the quasi-particle band in the \([HH]\) direction on surface B, with white dashed lines indicating the boundary of the integral (Eq. 1) used to calculate the RXS intensity. (B) Calculated RXS intensity, which exhibits a broad kink starting around \( H = 0.2 \text{ rlu} \). The momentum range, where the broad peak appears, is indicated by green dashed lines on both (A) and (B). (C) Experimentally measured RXS intensity on CeCoIn5 at \( T = 22 \text{ K} \), which shows a broad peak around \( 0.2 < H < 0.4 \text{ rlu} \). The corresponding momentum range is indicated by dashed lines.

**DISCUSSION**

Here, we showed the results of a complementary SI-STM and RXS study that probes the significance of band structure effects in RXS scattering measurements. The temperature, material, and photon energy dependences of the RXS data indicate their sensitivity to the formation of the Kondo lattice in the CeM1In5 system. On the basis of the experimental data, our model calculations show that the strongly dispersing \( f \)-bands can give rise to an enhancement of the RXS in a similar momentum range. These observations suggest that the bulk quasi-particle interference, as proposed by Abbamonte et al. (21) and Dalla Torre et al. (22),
is responsible for the RXS signal in the present measurements. We should emphasize that although the quasi-particle interference in STM conductance maps arises from native defects, atomic step edges, and impurities, it may have an additional contribution in RXS measurements. When an x-ray photon excites a core electron to the valence band, it creates a localized core-hole potential, which can act as a scattering center. Therefore, this perturbation can also lead to the scattering and interference of the itinerant quasi-particles (45).

It must also be noted that orbital degrees of freedom and their structure factors play a key role in RXS and SI-STM measurements, as in the case of the d-wave form factor of the charge order in the cuprates (12, 13). To obtain the relation between RXS and SI-STM expressed in Eq. 1, Abbamonte et al. (21) purposely disregarded the atomic orbital components of the electronic wave function. This is a particularly inappropriate approximation for f-electron systems, where the atomic multiplet structure is usually well defined in core-hole spectroscopies. This opens the possibility, for example, that the broad scattering enhancement in the RXS is related to how the polarized light couples to particular orbitals (39), which might only become available to the x-ray scattering process after band hybridization below T*. In our experiments, this scenario seems unlikely, given that all our measurements were done in σ-scattering geometry (that is, light polarization constantly in the a-b plane of the sample throughout the momentum scans, as Fig. 2B shows). This suggests that the momentum structure observed in our RXS data is likely related to the band structure of CeCoIn₅, as supported by the temperature-dependent measurements and their similarity to the SI-STM signal. Nevertheless, given the complexity of the RXS cross section, a comprehensive treatment of the scattering process, which takes into account the atomic multiplets in f-electron systems, would provide a more natural explanation to our results.

Our experiments, which demonstrate the relevance of quasi-particle interference in RXS, may also be relevant to the cuprates. In the past few years, a universal charge-order instability emerged as the most exciting progress in the study of high-temperature superconductivity in cuprates (5–20). The ubiquitous nature of this electronic phenomenon in the bulk of hole- and electron-doped cuprates came from RXS experiments on (Y,Nd)Ba₂Cu₃O₇, Bi-2201, Bi-2212, HgBa₂CuO₄+d, and Nd₂₋ₓCeₓCuO₄ (6–11). These measurements reveal an incommensurate scattering peak, with correlation lengths ranging from 20 to 75 Å, depending on the material and doping. Making a parallel between our RXS experiment on CeCoIn₅ and those on the cuprates suggests that the bulk quasi-particle interference features might be present in the RXS signal of the latter, perhaps even in coexistence with the charge-order peak. At this point, only further experiments can clarify the impact of quasi-particle interference to the RXS measurements in the cuprates.

Our experiments demonstrate that RXS can be a powerful momentum- and energy-resolved probe of the bulk band structure of materials, even in the absence of any ordering phenomena. These results not only pave the way to future RXS experiments on f-electron materials but also support the complementary relationship between RXS and STM measurements.

MATERIALS AND METHODS
Sample growth, STM, and RXS measurement technique
The single-crystal samples used for the measurements were grown from excess indium at the Los Alamos National Laboratory. Small, flat crystals were oriented along the crystallographic axes and cut into sizes suitable for STM and RXS measurements (0.5 to 2 mm × 0.5 to 2 mm × 0.2 mm). The samples used for the STM measurements were cleaved perpendicular to the c axis in ultrahigh vacuum at room temperature and immediately inserted into our homebuilt variable-temperature STM. Differential conductance measurements were performed using standard lock-in techniques, with bias applied to the sample.

The reported RXS experiments were performed at the REIXS (Resonant Elastic and Inelastic X-ray Scattering) beamline of the Canadian Light Source using a four-circle diffractometer in the temperature range of 22 to 200 K (46) and at the UE46_PGM-1 beamline of the Helmholtz-Zentrum Berlin at BESSY II (Berliner Elektronenspeicherring für Synchrotronstrahlung) with a two-circle diffractometer between 10 and 200 K. Reciprocal-space scans were acquired by rocking the sample angle (θ) at a fixed detector position (θdet = 170°). The samples were preoriented using Laue diffraction.

SUPPLEMENTARY MATERIALS
Supplementary material for this article is available at http://advances.sciencemag.org/cgi/content/full/2/10/e1601086/DC1

STM measurements and data analysis
RXS comparison on Cd-doped sample
L dependence of the RXS enhancement
Anomalous drop in the RXS signal
fig. 51. STM measurements on CeCoIn₅₋ₓCdₓ.  
fig. 52. Symmetrization of the conductance map. 
fig. 53. Comparison of the RXS cross section on CeCoIn₅₋ₓCdₓ. 
fig. 54. L dependence of the RXS measurement.

REFERENCES AND NOTES


Y. X. Wang, A. Chubukov, Charge-density-wave order with momentum (Q=0) and (Q,C) within the spin-fermion model: Continuous and discrete symmetry breaking, preemptive composite order, and relation to pseudogap in hole-doped cuprates. Phys. Rev. B 90, 035149 (2014).


P. Pfeiffer, Superconducting phases of Fe-electron compounds. Rev. Mod. Phys. 81, 1551–1624 (2009).


Acknowledgments: We thank R. Green and G. Sawatzky for stimulating discussions about the resonant x-ray process and Helmholtz-Zentrum Berlin for the allocation of synchrotron radiation beam time. Funding: The work at Princeton was primarily supported by a grant from the U.S. Department of Energy, Office of Basic Energy Sciences and Gordon and Betty Moore Foundation as part of EPiQS (Emergent Phenomena in Quantum Systems) Initiative (GBMF4530). The instrumentation and infrastructure at the Princeton Nanoscale Microscopy Laboratory used for this work were also supported by grants from NSF-DMR1104612, the NSF-MRSEC program through Princeton Center for Complex Materials (DMR-1420541), the Linda and Eric Schmidt Transformative Fund, and the W.M. Keck Foundation. The work at Binghamton University was supported by P.A.’s start-up funds. The work at Los Alamos was performed under the auspices of the U.S. Department of Energy, Office of Basic Energy Sciences, Division of Materials Sciences and Engineering. Part of the research described in this paper was performed at the Canadian Light Source, which is supported by the Canada Foundation for Innovation, Natural Sciences and Engineering Research Council of Canada (NSERC), the University of Saskatchewan, the Government of Saskatchewan, Western Economic Diversification Canada, the National Research Council Canada, and the Canadian Institutes of Health Research. The work at the University of British Columbia was supported by the Max Planck–University of British Columbia Centre for Quantum Materials; the Killam, Alfred P. Sloan, and NSERC’s Steacie Memorial Fellowships; the Alexander von Humboldt Fellowship; the Canada Research Chairs Program; and NSERC, Canada Foundation for Innovation, and the Canadian Institute for Advanced Research’s (CIFAR’s) Quantum Materials program. E.H.d.S.N. acknowledges support from the CIFAR Global Academy. Author contributions: A.G., P.A., and E.H.d.S.N. performed the STM measurements. E.H.d.S.N., A.G., P.A., R.S., E.S., F.H., and E.W. performed the RXS measurements. A.G., E.H.d.S.N., M.K., and P.A. analyzed the data. R.E.B., J.D.T., Z.F., and E.D.B. synthesized and characterized the materials. P.A., E.H.d.S.N., A.G., and A.Y. wrote the manuscript. All authors commented on the manuscript. Competing interests: The authors declare that they have no competing interests. Data and materials availability: All data needed to evaluate the conclusions in the paper are present in the paper and/or the Supplementary Materials. Additional data related to this paper may be requested from the authors.

Submitted 13 May 2016 Accepted 16 September 2016 Published 14 October 2016 10.1126/sciadv.1601086

Quasi-particle interference of heavy fermions in resonant x-ray scattering
doi: 10.1126/sciadv.1601086

This article is publisher under a Creative Commons license. The specific license under which this article is published is noted on the first page.

For articles published under CC BY licenses, you may freely distribute, adapt, or reuse the article, including for commercial purposes, provided you give proper attribution.

For articles published under CC BY-NC licenses, you may distribute, adapt, or reuse the article for non-commercial purposes. Commercial use requires prior permission from the American Association for the Advancement of Science (AAAS). You may request permission by clicking here.

The following resources related to this article are available online at http://advances.sciencemag.org. (This information is current as of May 30, 2017):

Updated information and services, including high-resolution figures, can be found in the online version of this article at:
http://advances.sciencemag.org/content/2/10/e1601086.full

Supporting Online Material can be found at:
http://advances.sciencemag.org/content/suppl/2016/10/11/2.10.e1601086.DC1

This article cites 45 articles, 8 of which you can access for free at:
http://advances.sciencemag.org/content/2/10/e1601086#BIBL